

## 5.2. Transformations of symmetry operations (motions)

BY H. ARNOLD

### 5.2.1. Transformations

Symmetry operations are transformations in which the coordinate system, *i.e.* the basis vectors **a**, **b**, **c** and the origin *O*, are considered to be at rest, whereas the object is mapped onto itself. This can be visualized as a ‘motion’ of an object in such a way that the object before and after the ‘motion’ cannot be distinguished.

A symmetry operation *W* transforms every point *X* with the coordinates *x, y, z* to a symmetrically equivalent point  $\tilde{X}$  with the coordinates  $\tilde{x}, \tilde{y}, \tilde{z}$ . In matrix notation, this transformation is performed by

$$\begin{pmatrix} \tilde{x} \\ \tilde{y} \\ \tilde{z} \end{pmatrix} = \begin{pmatrix} W_{11} & W_{12} & W_{13} \\ W_{21} & W_{22} & W_{23} \\ W_{31} & W_{32} & W_{33} \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} + \begin{pmatrix} w_1 \\ w_2 \\ w_3 \end{pmatrix} \\ = \begin{pmatrix} W_{11}x + W_{12}y + W_{13}z + w_1 \\ W_{21}x + W_{22}y + W_{23}z + w_2 \\ W_{31}x + W_{32}y + W_{33}z + w_3 \end{pmatrix}.$$

The  $(3 \times 3)$  matrix **W** is the rotation part and the  $(3 \times 1)$  column matrix **w** the translation part of the symmetry operation *W*. The pair  $(\mathbf{W}, \mathbf{w})$  characterizes the operation uniquely. Matrices **W** for point-group operations are given in Tables 11.2.2.1 and 11.2.2.2.

Again, we can introduce the augmented  $(4 \times 4)$  matrix (*cf.* Chapter 8.1)

$$\mathbb{W} = \begin{pmatrix} \mathbf{W} & \mathbf{w} \\ \mathbf{o} & \mathbf{1} \end{pmatrix} = \begin{pmatrix} W_{11} & W_{12} & W_{13} & w_1 \\ W_{21} & W_{22} & W_{23} & w_2 \\ W_{31} & W_{32} & W_{33} & w_3 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

The coordinates  $\tilde{x}, \tilde{y}, \tilde{z}$  of the point  $\tilde{X}$ , symmetrically equivalent to *X* with the coordinates *x, y, z*, are obtained by

$$\begin{pmatrix} \tilde{x} \\ \tilde{y} \\ \tilde{z} \\ 1 \end{pmatrix} = \begin{pmatrix} W_{11} & W_{12} & W_{13} & w_1 \\ W_{21} & W_{22} & W_{23} & w_2 \\ W_{31} & W_{32} & W_{33} & w_3 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \\ 1 \end{pmatrix} \\ = \begin{pmatrix} W_{11}x + W_{12}y + W_{13}z + w_1 \\ W_{21}x + W_{22}y + W_{23}z + w_2 \\ W_{31}x + W_{32}y + W_{33}z + w_3 \\ 1 \end{pmatrix},$$

or, in short notation,

$$\tilde{\mathbf{x}} = \mathbb{W}\mathbf{x}.$$

A sequence of symmetry operations can be obtained as a product of  $(4 \times 4)$  matrices  $\mathbb{W}$ .

An affine transformation of the coordinate system transforms the coordinates **x** of the starting point

$$\mathbf{x}' = \mathbb{Q}\mathbf{x}$$

as well as the coordinates  $\tilde{\mathbf{x}}$  of a symmetrically equivalent point

$$\begin{aligned} \tilde{\mathbf{x}}' &= \mathbb{Q}\tilde{\mathbf{x}} \\ &= \mathbb{Q}\mathbb{W}\mathbf{x} \\ &= \mathbb{Q}\mathbb{W}\mathbb{P}\mathbb{Q}\mathbf{x} \quad (\text{with } \mathbb{P} = \mathbb{Q}^{-1}) \\ &= \mathbb{Q}\mathbb{W}\mathbb{P}\mathbf{x}'. \end{aligned}$$

Thus, the affine transformation transforms also the symmetry-operation matrix  $\mathbb{W}$  and the new matrix  $\mathbb{W}'$  is obtained by

$$\mathbb{W}' = \mathbb{Q}\mathbb{W}\mathbb{P}.$$

*Example*

Space group *P4/n* (85) is listed in the space-group tables with two origins; origin choice 1 with  $\bar{4}$ , origin choice 2 with  $\bar{1}$  as point symmetry of the origin. How does the matrix  $\mathbb{W}$  of the symmetry operation  $\bar{4}^+ 0, 0, z; 0, 0, 0$  of origin choice 1 transform to the matrix  $\mathbb{W}'$  of symmetry operation  $\bar{4}^+ \frac{1}{4}, -\frac{1}{4}, z; \frac{1}{4}, -\frac{1}{4}, 0$  of origin choice 2?

In the space-group tables, origin choice 1, the transformed coordinates  $\tilde{x}, \tilde{y}, \tilde{z} = y, \bar{x}, \bar{z}$  are listed. The translation part is zero, *i.e.*  $\mathbf{w} = (0/0/0)$ . In Table 11.2.2.1, the matrix **W** can be found. Thus, the  $(4 \times 4)$  matrix  $\mathbb{W}$  is obtained:

$$\mathbb{W} = \begin{pmatrix} \mathbf{W} & \mathbf{w} \\ \mathbf{o} & \mathbf{1} \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ \bar{1} & 0 & 0 & 0 \\ 0 & 0 & \bar{1} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

The transformation to origin choice 2 is accomplished by a shift vector **p** with components  $\frac{1}{4}, -\frac{1}{4}, 0$ . Since this is a pure shift, the matrices **P** and **Q** are the unit matrix *I*. Now the shift vector **q** is derived:  $\mathbf{q} = -\mathbf{P}^{-1}\mathbf{p} = -\mathbf{I}\mathbf{p} = -\mathbf{p}$ . Thus, the matrices **P** and **Q** are

$$\mathbb{P} = \begin{pmatrix} 1 & 0 & 0 & \frac{1}{4} \\ 0 & 1 & 0 & \frac{1}{4} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \mathbb{Q} = \begin{pmatrix} 1 & 0 & 0 & \frac{1}{4} \\ 0 & 1 & 0 & \frac{1}{4} \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

By matrix multiplication, the new matrix  $\mathbb{W}'$  is obtained:

$$\mathbb{W}' = \mathbb{Q}\mathbb{W}\mathbb{P} = \begin{pmatrix} 0 & 1 & 0 & \frac{1}{2} \\ \bar{1} & 0 & 0 & 0 \\ 0 & 0 & \bar{1} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

If the matrix  $\mathbb{W}'$  is applied to  $x', y', z'$ , the coordinates of the starting point in the new coordinate system, we obtain the transformed coordinates  $\tilde{x}', \tilde{y}', \tilde{z}'$ ,

$$\begin{pmatrix} \tilde{x}' \\ \tilde{y}' \\ \tilde{z}' \\ 1 \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 & \frac{1}{2} \\ \bar{1} & 0 & 0 & 0 \\ 0 & 0 & \bar{1} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x' \\ y' \\ z' \\ 1 \end{pmatrix} = \begin{pmatrix} y' - \frac{1}{2} \\ \bar{x}' \\ \bar{z}' \\ 1 \end{pmatrix}.$$

By adding a lattice translation **a**, the transformed coordinates  $y + \frac{1}{2}, \bar{x}, \bar{z}$  are obtained as listed in the space-group tables for origin choice 2.

### 5.2.2. Invariants

A crystal structure and its physical properties are independent of the choice of the unit cell. This implies that invariants occur, *i.e.* quantities which have the same values before and after the transformation. Only some important invariants are considered in this section. Invariants of higher order (tensors) are treated by Altmann & Herzog (1994), second cumulant tensors, *i.e.* anisotropic temperature factors, are given in *International Tables for Crystallography* (2004), Vol. C.