

1. INTRODUCTION TO SPACE-GROUP SYMMETRY

Table 1.5.4.2

Additional symmetry operations due to a centring vector **t** and their locations

Symmetry operation at the origin		Additional symmetry operations									Representative space groups (numbers)
Symbol	Location	<i>C</i> , $t(\frac{1}{2}, \frac{1}{2}, 0)$		<i>A</i> , $t(0, \frac{1}{2}, \frac{1}{2})$		<i>B</i> , $t(\frac{1}{2}, 0, \frac{1}{2})$		<i>I</i> , $t(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$		<i>F</i>	
<i>m</i>	0, <i>y</i> , <i>z</i>	<i>b</i>	$\frac{1}{4}, y, z$	<i>n</i>	0, <i>y</i> , <i>z</i>	<i>c</i>	$\frac{1}{4}, y, z$	<i>n</i>	$\frac{1}{4}, y, z$	<i>b</i> , <i>n</i> , <i>c</i>	<i>Cmmm</i> , <i>Ammm</i> , <i>Bmmm</i> (65), <i>Immm</i> (71), <i>Fmmm</i> (69), <i>Cccm</i> , <i>Amaa</i> , <i>Bbmb</i> (66), <i>Ibca</i> (73), <i>Fddd</i> (70)
<i>c</i>		<i>n</i>		<i>b</i>		<i>m</i>		<i>b</i>			
<i>b</i>		<i>m</i>		<i>c</i>		<i>n</i>		<i>c</i>			
$d(0, \frac{1}{4}, \frac{1}{4})$		$d(0, \frac{3}{4}, \frac{1}{4})$		$d(0, \frac{3}{4}, \frac{3}{4})$		$d(0, \frac{1}{4}, \frac{3}{4})$				<i>d</i> , <i>d</i> , <i>d</i>	
<i>m</i>	<i>x</i> , 0, <i>z</i>	<i>a</i>	$x, \frac{1}{4}, z$	<i>c</i>	$x, \frac{1}{4}, z$	<i>n</i>	<i>x</i> , 0, <i>z</i>	<i>n</i>	$x, \frac{1}{4}, z$	<i>a</i> , <i>c</i> , <i>n</i>	As above
<i>a</i>		<i>m</i>		<i>n</i>		<i>c</i>		<i>c</i>			
<i>c</i>		<i>n</i>		<i>m</i>		<i>a</i>		<i>a</i>			
$d(\frac{1}{4}, 0, \frac{1}{4})$		$d(\frac{3}{4}, 0, \frac{1}{4})$		$d(\frac{1}{4}, 0, \frac{3}{4})$		$d(\frac{3}{4}, 0, \frac{3}{4})$				<i>d</i> , <i>d</i> , <i>d</i>	
<i>m</i>	<i>x</i> , <i>y</i> , 0	<i>n</i>	<i>x</i> , <i>y</i> , 0	<i>b</i>	$x, y, \frac{1}{4}$	<i>a</i>	$x, y, \frac{1}{4}$	<i>n</i>	$x, y, \frac{1}{4}$	<i>n</i> , <i>b</i> , <i>a</i>	As above
<i>b</i>		<i>a</i>		<i>m</i>		<i>n</i>		<i>a</i>			
<i>a</i>		<i>b</i>		<i>n</i>		<i>m</i>		<i>b</i>			
$d(\frac{1}{4}, \frac{1}{4}, 0)$		$d(\frac{3}{4}, \frac{3}{4}, 0)$		$d(\frac{1}{4}, \frac{3}{4}, 0)$		$d(\frac{3}{4}, \frac{1}{4}, 0)$				<i>d</i> , <i>d</i> , <i>d</i>	
<i>m</i>	<i>x</i> , <i>x</i> , <i>z</i>	$g(\frac{1}{2}, \frac{1}{2}, 0)$	<i>x</i> , <i>x</i> , <i>z</i>	$g(\frac{1}{4}, \frac{1}{4}, \frac{1}{2})$	$x, x + \frac{1}{4}, z$	$g(\frac{1}{4}, \frac{1}{4}, \frac{1}{2})$	$x, x - \frac{1}{4}, z$	$n(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$	<i>x</i> , <i>x</i> , <i>z</i>	<i>g</i> , <i>g</i> , <i>g</i>	<i>I4mm</i> (107), <i>F43m</i> (216), <i>F43c</i> (219), <i>I43d</i> (220)
<i>c</i>		$n(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$		$g(\frac{1}{4}, \frac{1}{4}, 0)$		$g(\frac{1}{4}, \frac{1}{4}, 0)$		$g(\frac{1}{2}, \frac{1}{2}, 0)$		<i>n</i> , <i>g</i> , <i>g</i>	
$d(\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$								$d(\frac{3}{4}, \frac{3}{4}, \frac{3}{4})$			
2	<i>x</i> , 0, 0	2 ₁	$x, \frac{1}{4}, 0$	2	$x, \frac{1}{4}, \frac{1}{4}$	2 ₁	$x, 0, \frac{1}{4}$	2 ₁	$x, \frac{1}{4}, \frac{1}{4}$	2 ₁ , 2, 2 ₁	
2	0, <i>y</i> , 0	2 ₁	$\frac{1}{4}, y, 0$	2 ₁	$0, y, \frac{1}{4}$	2	$\frac{1}{4}, y, \frac{1}{4}$	2 ₁	$\frac{1}{4}, y, \frac{1}{4}$	2 ₁ , 2 ₁ , 2	
2	0, 0, <i>z</i>	2	$\frac{1}{4}, \frac{1}{4}, z$	2 ₁	$0, \frac{1}{4}, z$	2 ₁	$\frac{1}{4}, 0, z$	2 ₁	$\frac{1}{4}, \frac{1}{4}, z$	2, 2 ₁ , 2 ₁	
2	<i>x</i> , \bar{x} , 0	2	$x, \bar{x} + \frac{1}{2}, 0$	2 ₁	$(-\frac{1}{4}, \frac{1}{4}, 0)$	2 ₁	$(-\frac{1}{4}, \frac{1}{4}, 0)$	2	$x, \bar{x}, \frac{1}{4}$	2, 2 ₁ , 2 ₁	
4	0, 0, <i>z</i>	4	$0, \frac{1}{2}, z$	4 ₂	$-\frac{1}{4}, \frac{1}{4}, z$	4 ₂	$\frac{1}{4}, \frac{1}{4}, z$	4 ₂	$0, \frac{1}{2}, z$	4, 4 ₂ , 4 ₂	
4 ₁	0, 0, <i>z</i>	4 ₁	$0, \frac{1}{2}, z$	4 ₃	$-\frac{1}{4}, \frac{1}{4}, z$	4 ₃	$\frac{1}{4}, \frac{1}{4}, z$	4 ₃	$0, \frac{1}{2}, z$	4 ₁ , 4 ₃ , 4 ₃	
$\bar{1}$	0, 0, 0	$\bar{1}$	$\frac{1}{4}, \frac{1}{4}, 0$	$\bar{1}$	$0, \frac{1}{4}, \frac{1}{4}$	$\bar{1}$	$\frac{1}{4}, 0, \frac{1}{4}$	$\bar{1}$	$\frac{1}{4}, \frac{1}{4}, \frac{1}{4}$	$\bar{1}, \bar{1}, \bar{1}$	

This example illustrates in particular the occurrence of symmetry elements of screw or glide type even in the case of symmorphic space groups where all coset representatives $W = (W, w)$ with respect to the translation subgroup can be chosen with $w = o$.

Note that, mainly for historical reasons, the screw rotations resulting from the threefold rotation along the [111] direction are not included in the extended Hermann–Mauguin symbol of cubic space groups, cf. Table 1.5.4.4. However, these screw rotations are represented in the cubic symmetry-element diagrams by the symbols



(cf. Table 2.1.2.7), as can be observed in the symmetry-element diagram for a group of type *P23* (195) in Fig. 1.5.4.1.

Example 4

A twofold rotation $W = y, x, \bar{z}$ with the line $x, x, 0$ as geometric element has linear part

$$W = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & \bar{1} \end{pmatrix}.$$

The composition W' of W with the translation $t(0, 1, 0)$ has

intrinsic translation part $w'_g = \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ 0 \end{pmatrix}$ and location part $w'_l = \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$. Since $(I - W)p = w'_l$ for $p = \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \end{pmatrix}$, the

symmetry operation $W' = y, x + \frac{1}{2}, \bar{z}$ is a screw rotation $2(\frac{1}{2}, \frac{1}{2}, 0)x, x + \frac{1}{2}, 0$ with the line $x, x + \frac{1}{2}, 0$ as geometric element and is thus of a different type to W (cf. Table 1.5.4.1).

In an *I*-centred lattice, the composition of W with the centring translation $t(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ has intrinsic translation part $w'_g = \begin{pmatrix} \frac{1}{2} \\ 0 \\ 0 \end{pmatrix}$ and location part $w'_l = \begin{pmatrix} 0 \\ 0 \\ \frac{1}{2} \end{pmatrix}$. One has $(I - W)p = w'_l$ for $p = \begin{pmatrix} 0 \\ 0 \\ \frac{1}{2} \end{pmatrix}$, hence the symmetry operation $W' = y + \frac{1}{2}, x + \frac{1}{2}, \bar{z} + \frac{1}{2}$ is a screw rotation $2(\frac{1}{2}, \frac{1}{2}, 0)x, x, \frac{1}{4}$ with the line $x, x, \frac{1}{4}$ as geometric element and is thus of a different type to W .

On the other hand, the translation subgroup T also contains the translation $t(\frac{1}{2}, -\frac{1}{2}, \frac{1}{2})$. In this case, the intrinsic translation part of $W' = y + \frac{1}{2}, x - \frac{1}{2}, \bar{z} + \frac{1}{2}$ is $w'_g = o$, hence W' is of the same type as W , i.e. a twofold rotation. The location part is

$w'_l = \begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$ and since $(I - W)p = w'_l$ for $p = \begin{pmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{4} \end{pmatrix}$, the geometric element of W' is the line $x + \frac{1}{2}, x, \frac{1}{4}$.