

## 5. DYNAMICAL THEORY AND ITS APPLICATIONS

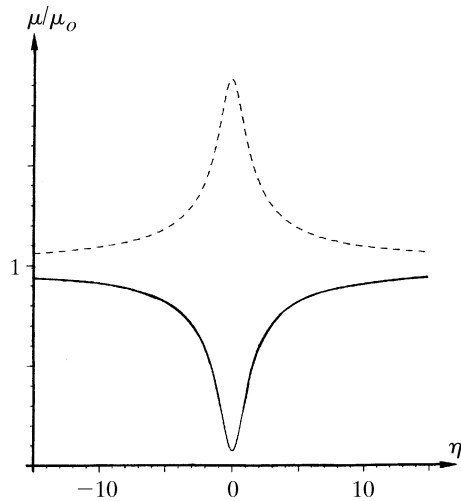


Fig. 5.1.6.1. Variation of the effective absorption with the deviation parameter in the transmission case for the 400 reflection of GaAs using Cu  $K\alpha$  radiation. Solid curve: branch 1; broken curve: branch 2.

$$\mu_j = \mu_o \left[ 1 \mp \frac{|C| |F_{ih}/F_{io}| \cos \varphi}{(\eta_r^2 + 1)^{1/2}} \right].$$

Fig. 5.1.6.1 shows the variations of the effective absorption coefficient  $\mu_j$  with  $\eta_r$  for wavefields belonging to branches 1 and 2 in the case of the 400 reflection of GaAs with Cu  $K\alpha$  radiation. It can be seen that for  $\eta_r = 0$  the absorption coefficient for branch 1 becomes significantly smaller than the normal absorption coefficient,  $\mu_o$ . The minimum absorption coefficient,  $\mu_o(1 - |CF_{ih}/F_{io}| \cos \varphi)$ , depends on the nature of the reflection through the structure factor and on the temperature through the Debye–Waller factor included in  $F_{ih}$  [equation (5.1.2.10b)] (Ohtsuki, 1964, 1965). For instance, in diamond-type structures, it is smaller for reflections with even indices than for reflections with odd indices. The influence of temperature is very important when  $|F_{ih}/F_{io}|$  is close to one; for example, for germanium 220 and Mo  $K\alpha$  radiation, the minimum absorption coefficient at 5 K is reduced to about 1% of its normal value,  $\mu_o$  (Ludewig, 1969).

### 5.1.6.2. Boundary conditions for the amplitudes at the entrance surface – intensities of the reflected and refracted waves

Let us consider an infinite plane wave incident on a crystal plane surface of infinite lateral extension. As has been shown in Section 5.1.3, two wavefields are excited in the crystal, with tie points  $P_1$  and  $P_2$ , and amplitudes  $D_{o1}, D_{h1}$  and  $D_{o2}, D_{h2}$ , respectively. Maxwell's boundary conditions (see Section A5.1.1.2 of the Appendix) imply continuity of the tangential component of the electric field and of the normal component of the electric displacement across the boundary. Because the index of refraction is so close to unity, one can assume to a very good approximation that there is continuity of the three components of both the electric field and the electric displacement. As a consequence, it can easily be shown that, *along the entrance surface*, for all components of the electric displacement

$$\begin{aligned} D_o^{(a)} &= D_{o1} + D_{o2} \\ 0 &= D_{h1} + D_{h2}, \end{aligned} \quad (5.1.6.2)$$

where  $D_o^{(a)}$  is the amplitude of the incident wave.

Using (5.1.3.11), (5.1.5.2) and (5.1.6.2), it can be shown that the intensities of the four waves are

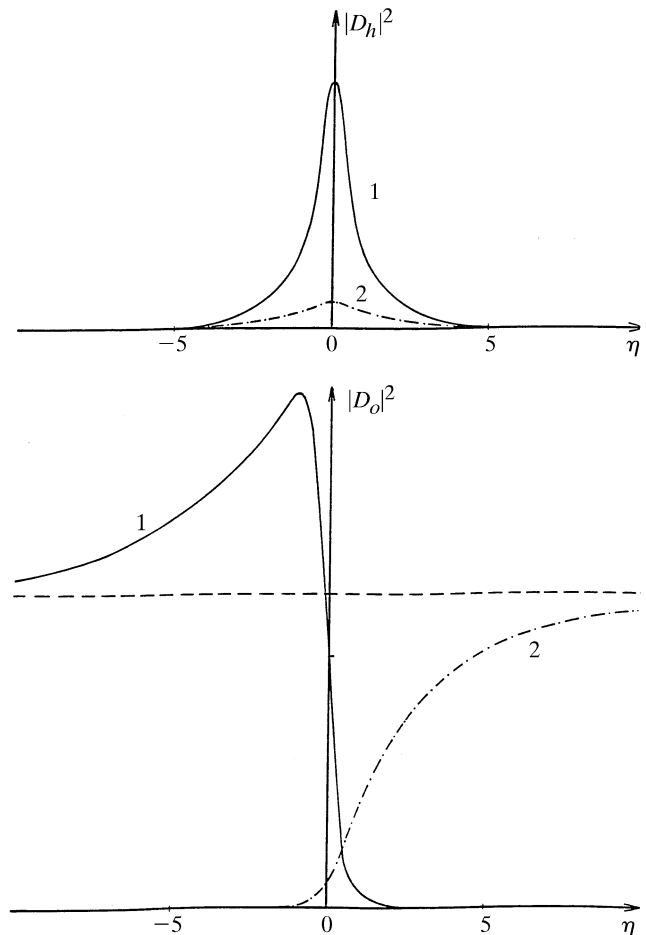


Fig. 5.1.6.2. Variation of the intensities of the reflected and refracted waves in an absorbing crystal for the 220 reflection of Si using Mo  $K\alpha$  radiation,  $t = 1$  mm ( $\mu t = 1.42$ ). Solid curve: branch 1; dashed curve: branch 2.

$$\begin{aligned} |D_{oj}|^2 &= |D_o^{(a)}|^2 \exp(-\mu_j z / \gamma_o) \left[ (1 + \eta_r^2)^{1/2} \mp \eta_r \right]^2 \\ &\quad \times [4(1 + \eta_r^2)]^{-1}, \end{aligned} \quad (5.1.6.3)$$

$$|D_{hj}|^2 = |D_o^{(a)}|^2 \exp(-\mu_j z / \gamma_o) |F_h / F_{\bar{h}}| [4\gamma(1 + \eta_r^2)]^{-1};$$

top sign:  $j = 1$ ; bottom sign:  $j = 2$ .

Fig. 5.1.6.2 represents the variations of these four intensities with the deviation parameter. Far from the reflection domain,  $|D_{h1}|^2$  and  $|D_{h2}|^2$  tend toward zero, as is normal, while

$$|D_{o1}|^2 \gg |D_{o2}|^2 \text{ for } \eta_r \Rightarrow -\infty,$$

$$|D_{o1}|^2 \ll |D_{o2}|^2 \text{ for } \eta_r \Rightarrow +\infty.$$

This result shows that the wavefield of highest intensity 'jumps' from one branch of the dispersion surface to the other across the reflection domain. This is an important property of dynamical theory which also holds in the Bragg case and when a wavefield crosses a highly distorted region in a deformed crystal [the so-called *interbranch scattering*; see, for instance, Authier & Balibar (1970) and Authier & Malgrange (1998)].

### 5.1.6.3. Boundary conditions at the exit surface

#### 5.1.6.3.1. Wavevectors

When a wavefield reaches the exit surface, it breaks up into its two constituent waves. Their wavevectors are obtained by applying