

1.1. INTRODUCTION TO THE PROPERTIES OF TENSORS

Furthermore, the left-hand term of (1.1.4.11) remains unchanged if we interchange the indices i and j . The terms on the right-hand side therefore also remain unchanged, whatever the value of T_{ll} or T_{kl} . It follows that

$$s_{ijll} = s_{jill}$$

$$s_{ijkl} = s_{ijlk} = s_{jikl} = s_{jilk}.$$

Similar relations hold for c_{ijkl} , Q_{ijkl} , p_{ijkl} and π_{ijkl} : the submatrices **2** and **3**, **4** and **7**, **5**, **6**, **8** and **9**, respectively, are equal.

Equation (1.4.1.11) can be rewritten, introducing the coefficients of the Voigt strain matrix:

$$S_\alpha = S_{ii} = \sum_l s_{iill} T_{ll} + \sum_{k \neq l} (s_{iikl} + s_{iilk}) T_{kl} \quad (\alpha = 1, 2, 3)$$

$$S_\alpha = S_{ij} + S_{ji} = \sum_l (s_{ijll} + s_{jill}) T_{ll} + \sum_{k \neq l} (s_{ijkl} + s_{ijlk} + s_{jikl} + s_{jilk}) T_{kl} \quad (\alpha = 4, 5, 6).$$

We shall now introduce a two-index notation for the elastic compliances, according to the following conventions:

$$\left. \begin{aligned} i = j; \quad k = l; \quad s_{\alpha\beta} &= s_{iill} \\ i = j; \quad k \neq l; \quad s_{\alpha\beta} &= s_{iikl} + s_{iilk} \\ i \neq j; \quad k = l; \quad s_{\alpha\beta} &= s_{ijkk} + s_{jikl} \\ i \neq j; \quad k \neq l; \quad s_{\alpha\beta} &= s_{ijkl} + s_{ijlk} + s_{jikl} + s_{jilk}. \end{aligned} \right\} \quad (1.1.4.12)$$

We have thus associated with the fourth-rank tensor a square 6×6 matrix with 36 coefficients:

β	1	2	3	4	5	6
α	11	12	13	14	15	16
2	21	22	23	24	25	26
3	31	32	33	34	35	36
4	41	42	43	44	45	46
5	51	52	53	54	55	56
6	61	62	63	64	65	66

One can translate relation (1.1.4.12) using the 9×9 matrix representing s_{ijkl} by adding term by term the coefficients of submatrices **2** and **3**, **4** and **7** and **5**, **6**, **8** and **9**, respectively:

$$\left(\begin{array}{c} 1 \\ 2+3 \end{array} \right) = \left(\begin{array}{c|c} 1 & 2 \\ \hline 4+7 & 5+6 \\ & +8+9 \end{array} \right) \times \left(\begin{array}{c} 1 \\ 2+3 \end{array} \right)$$

Using the two-index notation, equation (1.1.4.9) becomes

$$S_\alpha = s_{\alpha\beta} T_\beta. \quad (1.1.4.13)$$

A similar development can be applied to the other fourth-rank tensors π_{ijkl} , which will be replaced by 6×6 matrices with 36 coefficients, according to the following rules.

(i) *Elastic stiffnesses*, c_{ijkl} and *elasto-optic coefficients*, p_{ijkl} :

$$\left(\begin{array}{c} 1 \\ 2 \end{array} \right) = \left(\begin{array}{c|c} 1 & 2 \\ \hline 4 & 5 \end{array} \right) \times \left(\begin{array}{c} 1 \\ 2 \end{array} \right)$$

where

$$c_{\alpha\beta} = c_{ijkl}$$

$$p_{\alpha\beta} = p_{ijkl}.$$

(ii) *Piezo-optic coefficients*, π_{ijkl} :

$$\left(\begin{array}{c} 1 \\ 2 \end{array} \right) = \left(\begin{array}{c|c} 1 & 2+3 \\ \hline 4 & 5+6 \end{array} \right) \times \left(\begin{array}{c} 1 \\ 2 \end{array} \right)$$

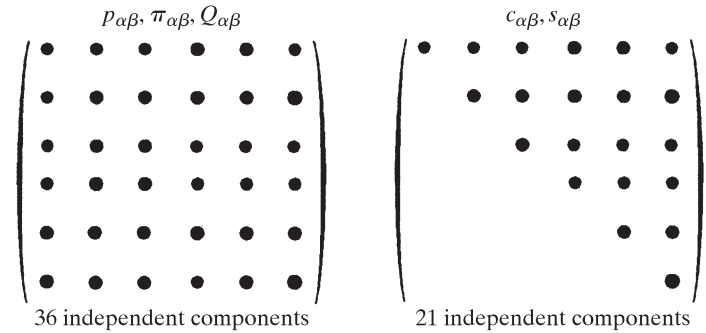
where

$$\left. \begin{aligned} i = j; \quad k = l; \quad \pi_{\alpha\beta} &= \pi_{iill} \\ i = j; \quad k \neq l; \quad \pi_{\alpha\beta} &= \pi_{iikl} + \pi_{iilk} \\ i \neq j; \quad k = l; \quad \pi_{\alpha\beta} &= \pi_{ijkk} = \pi_{jikl} \\ i \neq j; \quad k \neq l; \quad \pi_{\alpha\beta} &= \pi_{ijkl} + \pi_{jilk} = \pi_{ijlk} + \pi_{jikl}. \end{aligned} \right\}$$

(iii) *Electrostriction coefficients*, Q_{ijkl} : same relation as for the elastic compliances.

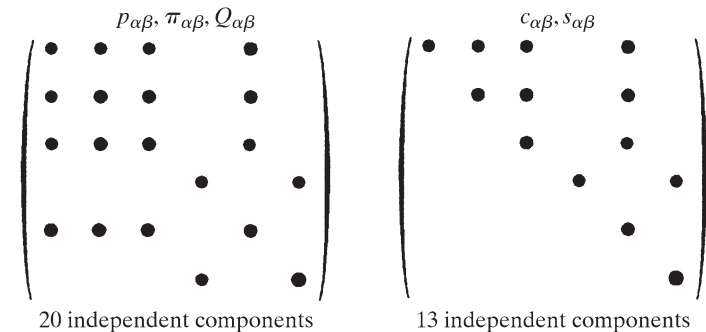
1.1.4.10.6. Independent components of the matrix associated with a fourth-rank tensor according to the following point groups

1.1.4.10.6.1. Triclinic system, groups $\bar{1}$, 1



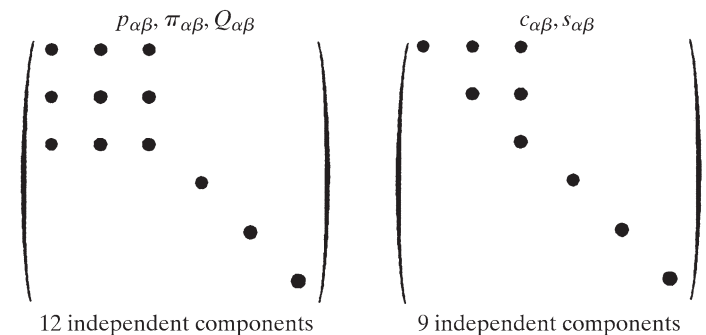
1.1.4.10.6.2. Monoclinic system

Groups $2/m$, 2 , m , twofold axis parallel to Ox_2 :



1.1.4.10.6.3. Orthorhombic system

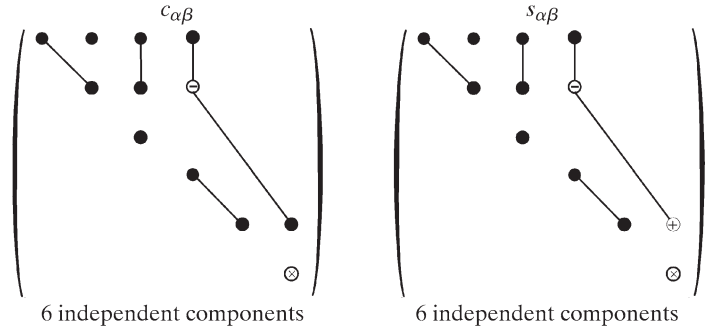
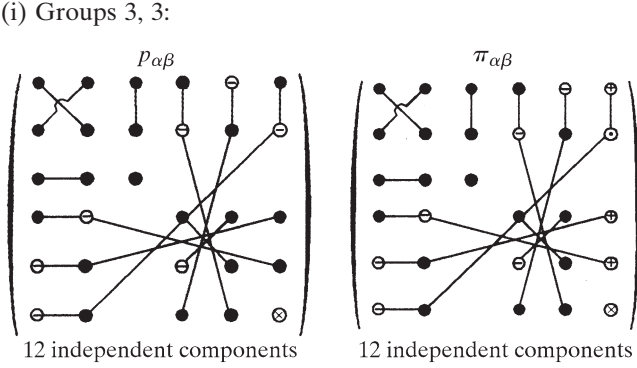
Groups mmm , $2mm$, 222 :



1. TENSORIAL ASPECTS OF PHYSICAL PROPERTIES

1.1.4.10.6.4. Trigonal system

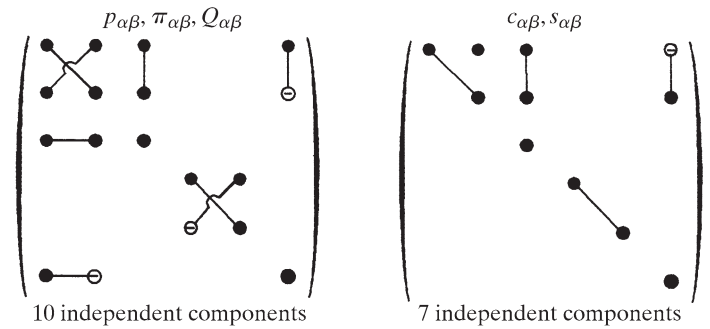
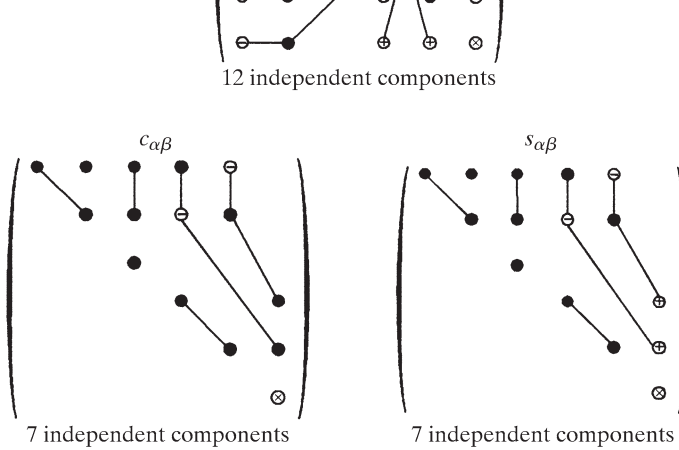
(i) Groups 3, $\bar{3}$:



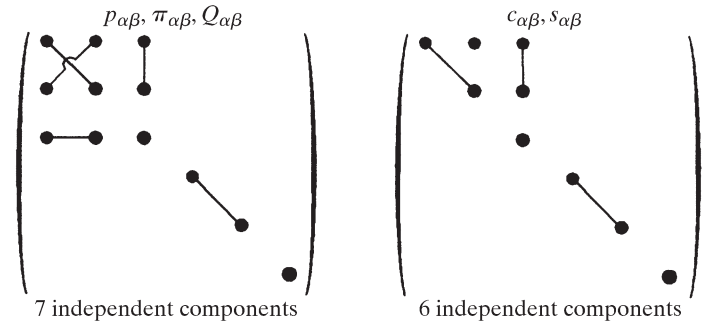
with the same conventions. The sign of c_{14} depends on the orientation of the Ox_1 axis.

1.1.4.10.6.5. Tetragonal system

(i) Groups 4, $\bar{4}$ and $4/m$:

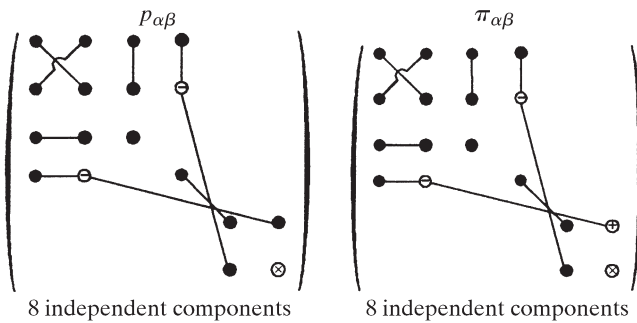


(ii) Groups $422, 4mm, \bar{4}2m$ and $4/m\bar{m}$:



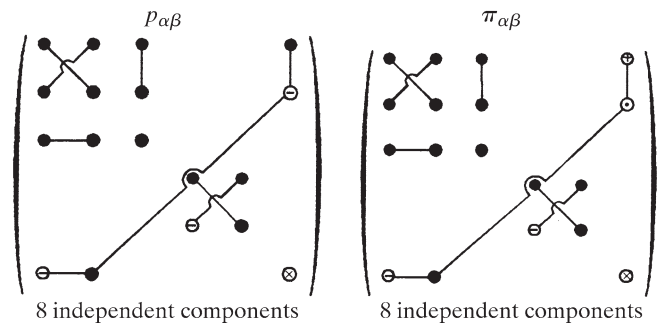
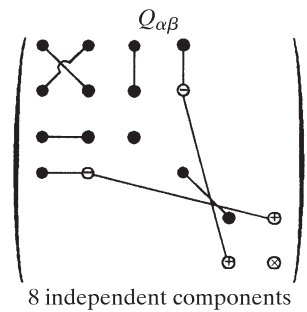
where \ominus is a component numerically equal but opposite in sign to the heavy dot component to which it is linked; \oplus is a component equal to twice the heavy dot component to which it is linked; \odot is a component equal to minus twice the heavy dot component to which it is linked; \otimes is equal to $1/2(p_{11} - p_{12})$, $(\pi_{11} - \pi_{12})$, $2(Q_{11} - Q_{12})$, $1/2(c_{11} - c_{12})$ and $2(s_{11} - s_{12})$, respectively.

(ii) Groups 32, $3m, \bar{3}m$:

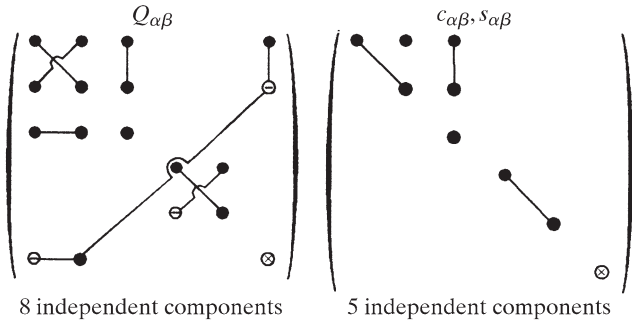


1.1.4.10.6.6. Hexagonal system

(i) Groups 6, $\bar{6}$ and $6/m$:

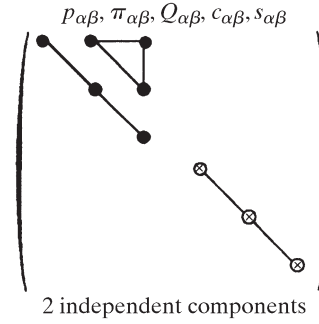


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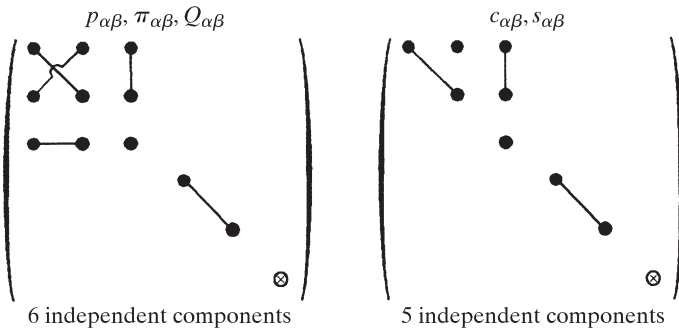


1.1.4.10.6.8. Spherical system

For all tensors



(ii) Groups 622, 6mm, $\bar{6}2m$ and 6/mmm:



1.1.4.10.7. Reduction of the number of independent components of axial tensors of rank 2

It was shown in Section 1.1.4.5.3.2 that axial tensors of rank 2 are actually tensors of rank 3 antisymmetric with respect to two indices. The matrix of independent components of a tensor such that

$$g_{ijk} = -g_{jik}$$

is given by

$$\left(\begin{array}{ccc|cc|cc} & 122 & 133 & 123 & 131 & 132 & 121 \\ -121 & & 223 & & 231 & -122 & 232 & -123 \\ -131 & -232 & & -233 & -132 & & -133 & -231 \end{array} \right).$$

The second-rank axial tensor g_{kl} associated with this tensor is defined by

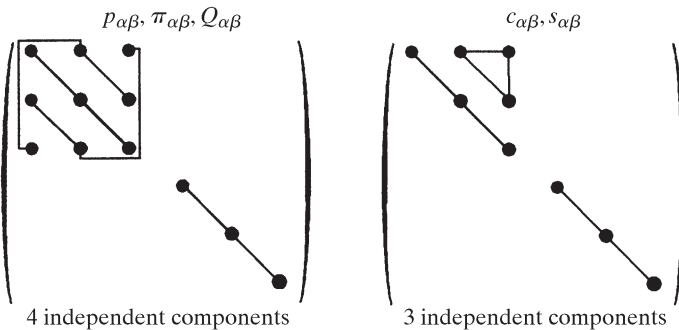
$$g_{kl} = \frac{1}{2}\epsilon_{ijk}g_{ijl}.$$

For instance, the piezomagnetic coefficients that give the magnetic moment M_i due to an applied stress T_α are the components of a second-rank axial tensor, $\Lambda_{i\alpha}$ (see Section 1.5.7.1):

$$M_i = \Lambda_{i\alpha}T_\alpha.$$

1.1.4.10.6.7. Cubic system

(i) Groups 23 and 3m:



1.1.4.10.7.1. Independent components according to the following point groups

(i) Triclinic system

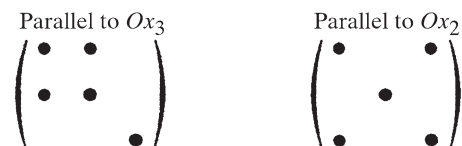
(a) Group 1:



(b) Group $\bar{1}$: all components are equal to zero.

(ii) Monoclinic system

(a) Group 2:



(ii) Groups 432, $\bar{4}3m$ and $m\bar{3}m$:

