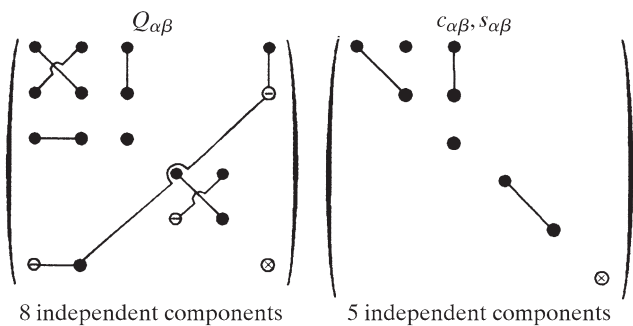
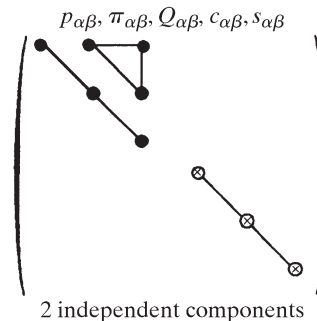


1.1. INTRODUCTION TO THE PROPERTIES OF TENSORS

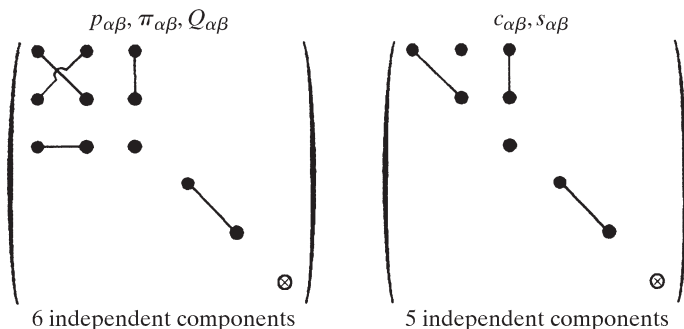


1.1.4.10.6.8. Spherical system

For all tensors



(ii) Groups 622, 6mm, 6̄2m and 6/mmm:



1.1.4.10.7. Reduction of the number of independent components of axial tensors of rank 2

It was shown in Section 1.1.4.5.3.2 that axial tensors of rank 2 are actually tensors of rank 3 antisymmetric with respect to two indices. The matrix of independent components of a tensor such that

$$g_{ijk} = -g_{jik}$$

is given by

$$\begin{pmatrix} & 122 & 133 & | & 123 & 131 & | & 132 & & 121 \\ -121 & & 223 & | & & 231 & -122 & | & 232 & -123 \\ -131 & -232 & & | & -233 & & -132 & | & & -133 & -231 \end{pmatrix}$$

The second-rank axial tensor  $g_{kl}$  associated with this tensor is defined by

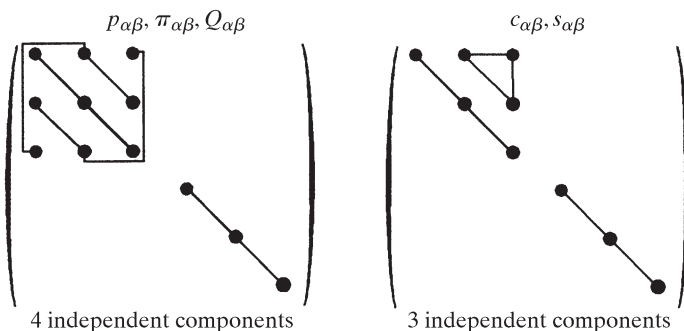
$$g_{kl} = \frac{1}{2} \epsilon_{ijk} g_{ijl}$$

For instance, the piezomagnetic coefficients that give the magnetic moment  $M_i$  due to an applied stress  $T_\alpha$  are the components of a second-rank axial tensor,  $\Lambda_{i\alpha}$  (see Section 1.5.7.1):

$$M_i = \Lambda_{i\alpha} T_\alpha$$

1.1.4.10.6.7. Cubic system

(i) Groups 23 and 3m:



1.1.4.10.7.1. Independent components according to the following point groups

(i) Triclinic system

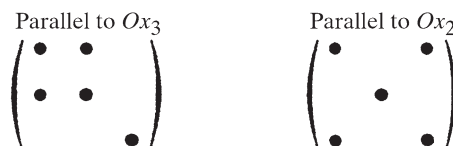
(a) Group 1:



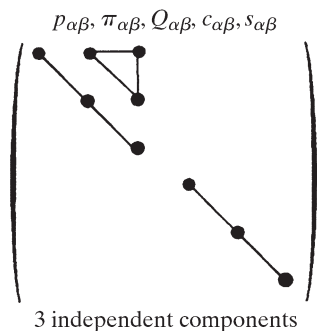
(b) Group 1̄: all components are equal to zero.

(ii) Monoclinic system

(a) Group 2:

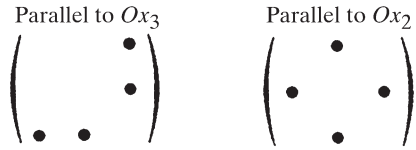


(ii) Groups 432, 4̄3m and m3̄m:



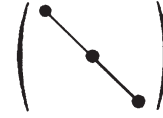
1. TENSORIAL ASPECTS OF PHYSICAL PROPERTIES

(b) Group  $m$ :



(v) Cubic and spherical systems

(a) Groups 23, 432 and  $\infty A_\infty$ :



(c) Group  $2/m$ : all components are equal to zero.

(iii) Orthorhombic system

(a) Group 222:



The axial tensor is reduced to a pseudoscalar.

(b) Groups  $m\bar{3}$ ,  $43m$ ,  $m\bar{3}m$  and  $\infty(A_\infty/M)C$ : all components are equal to zero.

1.1.4.10.7.2. Independent components of symmetric axial tensors according to the following point groups

Some axial tensors are also symmetric. For instance, the optical rotatory power of a gyrotropic crystal in a given direction of direction cosines  $\alpha_1, \alpha_2, \alpha_3$  is proportional to a quantity  $G$  defined by (see Section 1.6.5.4)

$$G = g_{ij}\alpha_i\alpha_j,$$

where the gyration tensor  $g_{ij}$  is an axial tensor. This expression shows that only the symmetric part of  $g_{ij}$  is relevant. This leads to a further reduction of the number of independent components:

(i) Triclinic system

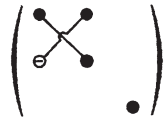
(a) Group 1:



(c) Group  $mmm$ : all components are equal to zero.

(iv) Trigonal, tetragonal, hexagonal and cylindrical systems

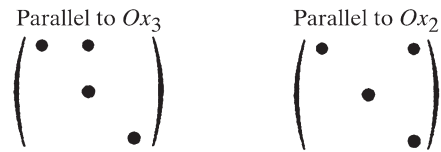
(a) Groups 3, 4, 6 and  $A_\infty$ :



(b) Group  $\bar{1}$ : all components are equal to zero.

(ii) Monoclinic system

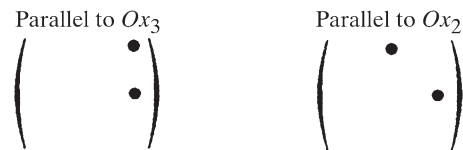
(a) Group 2:



(b) Groups 32, 42, 62 and  $A_\infty \infty A_2$ :



(b) Group  $m$ :



(c) Groups  $3m$ ,  $4m$ ,  $6m$  and  $A_\infty \infty M$ :



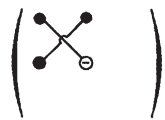
(c) Group  $2/m$ : all components are equal to zero.

(iii) Orthorhombic system

(a) Group 222:



(d) Group  $\bar{4}$ :



(e) Group  $\bar{4}2m$ :



(b) Group  $mm2$ :



(f) Groups  $\bar{3}$ ,  $4/m$ ,  $\bar{6}2m$ ,  $\bar{3}m$ ,  $4/m\bar{m}$  and  $6/m\bar{m}$ : all components are equal to zero.

(c) Group  $mmm$ : all components are equal to zero.