

1. TENSORIAL ASPECTS OF PHYSICAL PROPERTIES

$$P\Omega Q Q^{-1}v = DQ^{-1}v = 0,$$

where  $D_{ij} = d_i\delta_{ij}$ . The dimension of the solution space is the number of elements  $d_i$  that are equal to zero. The corresponding rows of  $Q$  form a basis for the solution space. (See example further on.)

1.2.7.4.4. Determination of the vector representation

For a point group  $G$ , its isomorphism class and its character table are known. For each conjugacy class, a representative element is given as word  $A_1A_2 \dots$  where the  $A_i$ 's correspond to generators. Replacing the letters by the generating matrices, one obtains as product a matrix for which the trace is the character of the vector representation in the conjugacy class. The characters of all conjugacy classes being known, the representation can be decomposed into irreducible components by means of

$$m_\alpha = (1/|G|) \sum_i n_i \chi_\alpha^*(i) \chi(i),$$

where  $\alpha$  labels the irreducible representations (the row number in the character table),  $m_\alpha$  the number of times the representation  $\alpha$  occurs,  $|G|$  the order of the group  $G$ ,  $n_i$  the number of elements in the  $i$ th conjugacy class (given as the second row in the character table),  $\chi_\alpha(i)$  the cyclotomic in the  $i$ th row and  $\alpha$ th column of the character table, and  $\chi(i)$  the calculated character in the  $i$ th conjugacy class.

1.2.7.4.5. Determination of tensor products and their decomposition

Given a character (for an irreducible representation from the character table, or for the vector representation, for example), the character of the standard rank  $n$  tensor is the  $n$ th power of the character and can be decomposed with the multiplicity formula for  $m_\alpha$  given above.

Fully symmetrized or antisymmetrized tensor products have characters given by

$$\begin{aligned} n = 2 : \chi^\pm(R) &= \frac{1}{2!} (\chi(R)^2 \pm \chi(R^2)) \\ n = 3 : \chi^\pm(R) &= \frac{1}{3!} (\chi(R)^3 \pm 3\chi(R^2)\chi(R) + 2\chi(R^3)) \\ n = 4 : \chi^\pm(R) &= \frac{1}{4!} (\chi(R)^4 \pm 6\chi(R^2)\chi(R)^2 + 3\chi(R^2)^2 \\ &\quad + 8\chi(R^3)\chi(R) \pm 6\chi(R^4)) \\ n = 5 : \chi^\pm(R) &= \frac{1}{5!} (\chi(R)^5 \pm 10\chi(R^2)\chi(R)^3 + 15\chi(R^2)^2\chi(R) \\ &\quad + 20\chi(R^3)\chi(R)^2 \pm 20\chi(R^3)\chi(R^2) \\ &\quad \pm 30\chi(R^4)\chi(R) + 24\chi(R^5)) \\ n = 6 : \chi^\pm(R) &= \frac{1}{6!} (\chi(R)^6 \pm 15\chi(R^2)\chi(R)^4 + 45\chi(R^2)^2\chi(R)^2 \\ &\quad + 40\chi(R^3)^2 \pm 15\chi(R^2)^3 + 40\chi(R^3)\chi(R)^3 \\ &\quad \pm 120\chi(R^3)\chi(R^2)\chi(R) \pm 90\chi(R^4)\chi(R)^2 \\ &\quad + 90\chi(R^4)\chi(R^2) + 144\chi(R^5)\chi(R) \\ &\quad \pm 120\chi(R^6)). \end{aligned}$$

From this follows immediately the dimension of the subspaces of symmetric and antisymmetric tensors:

$$\begin{aligned} n = 2 : & \frac{1}{2}(d^2 \pm d) \\ n = 3 : & \frac{1}{6}(d^3 \pm 3d^2 + 2d) \\ n = 4 : & \frac{1}{24}(d^4 \pm 6d^3 + 11d^2 \pm 6d) \\ n = 5 : & \frac{1}{120}(d^5 \pm 10d^4 + 35d^3 \pm 50d^2 + 24d) \\ n = 6 : & \frac{1}{720}(d^6 \pm 15d^5 + 85d^4 \pm 225d^3 + 274d^2 \pm 120d). \end{aligned}$$

The general expression for arbitrary rank can be determined as follows. (See also Section 1.2.2.7)

(1) If  $n$  is the rank, the first step is to determine all possible decompositions

$$n = \sum_{i=1}^n f_i$$

with non-negative integers  $f_i$  satisfying  $f_i \leq f_{i-1}$ .

(2) For each such decomposition  $m = 1, \dots, n_{\text{tot}}$  there is a term

$$P_m = \prod_{i=1}^p \binom{N_i}{f_i} (f_i - 1)!,$$

where  $N_1 = n$ ,  $N_i = N_{i-1} - f_{i-1}$  ( $i > 1$ ) and  $p$  is the number of nonzero integers  $f_i$ .

(3) If there are equal values of  $f_i$  in the  $m$ th decomposition,  $P_m$  should be divided by  $t!$  for each  $t$ -tuple of equal values ( $f_{k+1} = \dots = f_{k+t}$ ).

(4) The sign of the term  $P_m$  is  $+1$  for a symmetrized power and

$$\prod_{i=1}^p (-1)^{(f_i-1)}$$

for an antisymmetrized power.

(5) The expression for the character of the (anti)symmetrized power then is

$$\chi^\pm(R) = (1/M!) \sum_{m=1}^{n_{\text{tot}}} \text{sign}_m P_m \prod_{i=1}^p \chi(R^{f_i}).$$

1.2.7.4.6. Invariant tensors

Once one has the character of the properly symmetrized tensor, the number of invariants is just  $m_1$ , the number of times the trivial representation occurs in the decomposition.

Example (1). Dimension 3, rank 3, symmetry type (123), group 3. Basis:  $xxx, xxy, xxz, xyy, xyz, xzz, yyy, yyz, yzz, zzz$ . Under

$$\begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$$

the basis elements go to  $yyy, yyz, yyx, yzz, yzx, yxx, zzz, zzx, zxx, xxx$ , respectively, and these are equivalent to  $yyy, yyz, xyy, yzz, xyz, xxy, zzz, xzz, xxz, xxx$ , respectively. This gives the ten-dimensional matrix

1.2. REPRESENTATIONS OF CRYSTALLOGRAPHIC GROUPS

$$M = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \end{pmatrix}.$$

Then  $P(M - E)Q = D$ , with  $D$  diagonal. There are four diagonal elements of  $D$  which are zero, and the invariant tensors correspond to the corresponding four columns of the matrix  $Q$ . The invariant polynomials are

$$xxx + yyy + zzz, \quad xxy + xzz + yyz, \quad xxz + yzz + xyy, \quad xyz.$$

*Example (2).* Dimension 2, rank 2, symmetry type (12). Group generated by

$$\begin{pmatrix} 0 & -1 \\ 1 & -1 \end{pmatrix}.$$

Basis  $xx, xy, yy$  goes to  $yy, -xy + yy, xx - 2xy + yy$ . This gives

$$M = \begin{pmatrix} 0 & 0 & 1 \\ 0 & -1 & -2 \\ 1 & 1 & 1 \end{pmatrix}.$$

Because

$$\begin{pmatrix} 1 & 0 & 0 \\ 2 & 1 & 2 \\ 1 & 0 & 1 \end{pmatrix} (M - E) \begin{pmatrix} 1 & -1 & 1 \\ 0 & 1 & 0 \\ 0 & -1 & 1 \end{pmatrix} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 2 \end{pmatrix},$$

the invariant tensor corresponds to the second column of  $Q$ , which as a polynomial reads  $-xx + xy - yy$ . This can be written with the tensor  $T_{ij}$  as

$$-xx + xy - yy = - \sum_{i,j} T_{ij} x_i x_j, \quad T_{ij} = \begin{pmatrix} 1 & -\frac{1}{2} \\ -\frac{1}{2} & 1 \end{pmatrix}.$$

This tensor  $T$  is invariant under the group.

*Example (3).* Dimension 3, rank 2, tensor type (12). Group generated by matrix  $([[[0 -1 0][1 0 0][0 0 1]])$ . The basis  $xx, xy, xz, yy, yz, zz$  goes under the generator to  $yy, -xy, -yz, xx, xz, zz$ . The solution of  $(M - E)v = 0$  is

$$\alpha_1(xx + yy) + \alpha_2zz.$$

The matrix  $D$  has two zeros on the diagonal.

*Example (4).* Dimension 3, rank 3, type (123). Same group as in Example (3). Basis  $xxx, xxy, xxz, xyy, xyz, xzz, yyy, yyz, yzz, zzz$ . The solution

$$\alpha_1(xxz + yyz) + \alpha_2zzz$$

corresponds to a tensor with relations  $T_{113} = T_{223}, T_{111} = T_{112} = T_{122} = T_{123} = T_{133} = T_{222} = T_{233} = 0$ .

*Example (5).* Dimension 3, rank 4, type  $((12)(34))$ . Not only  $i_1 \leq i_2$  and  $i_3 \leq i_4$ , but also  $(i_1 i_2)$ , should come lexicographically before  $(i_3 i_4)$ . Basis  $xxxx, xxxy, xxxz, xxyy, xxyz, xxzz, xyxy, xyxz, xyyy, xyyz, xyzx, xzxx, xzyy, xzyz, xzzz, yyyy, yyyz, yyyz, yzyz, yzzz, zzzz$ . Under the same group as in example (3), there are seven invariants. Invariant polynomial:

$$\alpha_1(xxxx + yyyy) + \alpha_2(xxxy - xyxy) + \alpha_3(xxyy) + \alpha_4(xyxy) + \alpha_5(zzzz) + \alpha_6(xxzz + yyzz) + \alpha_7(xzxx + yzyz).$$

Table 1.2.7.2. Calculation with characters

Generator	Composite character	Characters				Decomposition
$\begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$ Example (1)	$R$	$E$	$A$	$AA$		
	$\chi(R)$	3	0	0		
	$\chi(R)^3$	27	0	0		
	$\chi(R^2)$	3	0	0		
	$\chi(R^2)\chi(R)$	9	0	0		
	$\frac{1}{6}(\chi(R)^3 + 3\chi(R^2)\chi(R) + 2\chi(R^3))$	10	1	1	$4D_1 + 3D_2 + 3D_3$	
$\begin{pmatrix} 0 & -1 \\ 1 & -1 \end{pmatrix}$ Example (2)	$R$	$E$	$A$	$AA$		
	$\chi(R)$	2	-1	-1		
	$\chi(R)^2$	4	1	1		
	$\chi(R^2)$	2	-1	-1		
	$\frac{1}{2}(\chi(R)^2 + \chi(R^2))$	3	0	0	$D_1 + D_2 + D_3$	
$\begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$ Example (3)	$R$	$E$	$A$	$AA$	$AAA$	
	$\chi(R)$	3	1	-1	1	
	$\chi(R)^2$	9	1	1	1	
	$\chi(R^2)$	3	-1	3	-1	
	$\frac{1}{2}(\chi(R)^2 + \chi(R^2))$	6	0	2	0	$2D_1 + D_2 + 2D_3 + D_4$
As above Example (4)	$\chi(R)$	3	1	-1	1	
	$\chi(R)^3$	27	1	-1	1	
	$\chi(R^2)$	3	-1	3	-1	
	$\chi(R^2)\chi(R)$	9	-1	-3	-1	
	$\chi(R^3)$	3	1	-1	1	
	$\frac{1}{6}(\chi(R)^3 + 3\chi(R^2)\chi(R) + 2\chi(R^3))$	10	0	-2	0	$2D_1 + 3D_2 + 2D_3 + 3D_4$
As above Example (5)	$\chi(R)$	3	1	-1	1	
	$\frac{1}{2}(\chi(R)^2 + \chi(R^2)) = \chi_s(R)$	6	0	2	0	
	$\chi_s(R)^2$	36	0	4	0	
	$\chi_s(R^2)$	6	2	6	2	
	$((12)(34))$	21	1	5	1	$7D_1 + 4D_2 + 6D_3 + 4D_4$
As above, example (6)	$\frac{1}{2}(\chi(R)^2 - \chi(R^2))$	3	1	-1	1	$D_1 + D_2 + D_4$
As above, example (7)	$\frac{1}{6}(\chi(R)^3 - 3\chi(R^2)\chi(R) + 2\chi(R^3))$	1	1	1	1	$D_1$

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This corresponds to the tensor relations

$$\begin{aligned}
 T_{xxxx} &= -T_{yyyy} & T_{xxyy} &= T_{yyxx} & T_{xxzz} &= 0 \\
 T_{xxyz} &= 0 & T_{xxzz} &= T_{yyzz} & T_{xyxz} &= 0 \\
 T_{xyyz} &= 0 & T_{xyzz} &= 0 & T_{xzxx} &= T_{yzzy} \\
 T_{xzyy} &= 0 & T_{xzyz} &= 0 & T_{xzzz} &= 0 \\
 T_{yyyz} &= 0 & T_{yzzz} &= 0 & & \\
 \end{aligned}$$

$$\rightarrow \begin{pmatrix} \alpha_1 & \alpha_3 & \alpha_6 & 0 & 0 & \alpha_2 \\ \alpha_3 & \alpha_1 & \alpha_6 & 0 & 0 & -\alpha_2 \\ \alpha_6 & \alpha_6 & \alpha_5 & 0 & 0 & 0 \\ 0 & 0 & 0 & \alpha_7 & 0 & 0 \\ 0 & 0 & 0 & 0 & \alpha_7 & 0 \\ \alpha_2 & -\alpha_2 & 0 & 0 & 0 & \alpha_4 \end{pmatrix}.$$

The latter form is that of an elastic tensor with the usual convention  $1 = xx, 2 = yy, 3 = zz, 4 = yz, 5 = xz, 6 = xy$ .

*Example (6).* Dimension 3, rank 2, type [12]. The same group as in example (3). Basis  $xy, xz, yz \rightarrow -yx, -yz, xz$ , which are equivalent to  $xy, -yz, xz$ . The transformation in the tensor space is

$$M = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \rightarrow \begin{pmatrix} 0 & 0 & 0 \\ 0 & -1 & 1 \\ 0 & -1 & -1 \end{pmatrix} v = 0:$$

$$v = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \sim xy.$$

There is just one invariant antisymmetric polynomial  $xy = -yx$  corresponding to the tensor

$$T = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}.$$

*Example (7).* Dimension 3, rank 3, type [123]. Basis  $xyz$  invariant under the group:  $xyz \rightarrow -yxz \sim xyz$ . The corresponding tensor is the fully antisymmetric rank 3 tensor:  $T_{ijk} = 1$  if  $ijk$  is an even permutation of 123,  $= -1$  if  $ijk$  is an odd permutation, and  $= 0$  if two or three indices are equal (permutation tensor, see Section 1.1.3.7.2).

*Example (8).* Calculation with characters. See Table 1.2.7.2.

*Example (9).* The action matrix for a pseudotensor.

Take the group  $4/m$  with generators

$$\begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

Consider the rank 3 pseudotensor (123). The action matrix is determined from the action of the generators  $A$  and  $B$  on the basis:

	$A$	$B$
$xxx$	$-yyy$	$-xxx$
$xyx$	$xyy$	$-xxy$
$xxz$	$yyz$	$xxz$
$xyy$	$-xxy$	$-xyy$
$xyz$	$-xyz$	$xyz$
$xzz$	$-yzz$	$-xzz$
$yyy$	$xxx$	$-yyy$
$yyz$	$xxz$	$yyz$
$yzz$	$xzz$	$-yzz$
$zzz$	$zzz$	$zzz$

Therefore, the action matrix becomes

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}.$$

After diagonalization, one finds two nonzero elements on the diagonal:

$$\begin{aligned}
 zzz &= a; & xxz &= yyz = b; \\
 xxx &= xxy = xyy = xyz = xzz = yyy = yzz = 0.
 \end{aligned}$$

## 1.2.8. Glossary

$T_{i_1 \dots i_n}$	tensor of rank $n$
$O(n)$	orthogonal group
$\mathbb{Z}$	ring of integers
$\mathbf{e}_i$	basis vectors
$g$	metric tensor
$K$	point group
$R$	orthogonal transformation
$C_m$	cyclic group of order $m$
$SO(n)$	special orthogonal group
$\mathbb{Z}^+$	positive integers
$D_n$	dihedral group of order $n$
$E$	unit transformation, matrix or element
$I$	inversion
$D(K)$	representation of $K$
$\Gamma(K)$	matrix representation of $K$
$ K $	order of $K$
$\oplus$	sum of spaces or operators
$\otimes$	tensor product
$\in$	element of
$\mathbf{a}$	basis of space or lattice